

## Introduction

### Problem considered:

Many-body system  $S$  with  $N$  constituents each with  $d$  internal levels: typical model for cold atoms in optical lattices. The coherent evolution is described by  $U = \exp[-iHt]$  ( $\hbar = 1$ ) with  $H$  the Hamiltonian. Computational cost of a direct evaluation of  $U$  scales as  $d^{3N}$ .

### Walk-sum:

Walk-sums approximate **analytically** any chosen **pieces** of  $U$  without any prior knowledge about the system under study. Computational cost at fixed order scales polynomially with  $N$ .

- Walk-sum express any quantum-evolution as a sum of walks on a graph
- Reduces many-body physics to few-body physics exactly
- Yields Feynman path-integrals for systems with continuous degrees of freedom
- Is part of a wider approach to matrix-functions valid independently of physics

## Walk-sum: physical derivation

Consider  $S$  to comprise two subsystems: **small  $S$**  with  $d_S$  levels and **the rest  $S'$**

→ **Idea:** if  $S'$  is frozen in configuration  $\mu$ , evolving  $S$  is **easy**

→ **Goal:** describe the evolution of  $S$  in terms of situations where **only  $S$**  evolves

- Consider the projectors on a complete orthogonal set of configurations of  $S'$ :

$$\hat{\epsilon}_\mu = \underbrace{|\mu\rangle\langle\mu|}_{\text{Projects } S' \text{ in configuration } \mu} \otimes \underbrace{\mathcal{I}_S}_{\text{Identity on } S} \quad \text{with} \quad \sum_\mu \hat{\epsilon}_\mu = \mathcal{I}_S$$

- "Freeze"  $S'$  in a configuration  $\mu$  by projecting  $S'$  onto  $\mu$  every  $\delta t$

$$\text{Expand } U(t) = \lim_{\delta t \rightarrow 0} \prod_{n=0}^{t/\delta t} \left( \sum_\mu \hat{\epsilon}_\mu \right) \delta U \quad \dots$$

$$\dots \text{ yields terms like } \hat{\epsilon}_\nu U(t) \hat{\epsilon}_\mu = \lim_{\delta t \rightarrow 0} \underbrace{\hat{\epsilon}_\nu \delta U \hat{\epsilon}_\nu \delta U \dots \hat{\epsilon}_\nu \delta U}_{S' \text{ jumps from } \mu \text{ to } \nu} \underbrace{\hat{\epsilon}_\nu \delta U \hat{\epsilon}_\mu}_{\text{Zeno measurement of } \hat{\epsilon}_\nu \Rightarrow S' \text{ frozen in } \nu} + \dots$$

S evolves according to a **small effective  $U_\nu$**

- General expansion has jump  $\mu \rightarrow \nu$  at any time, jumps to intermediary configurations  $\eta$  etc. We find

### Walk-sum formula

$$\hat{\epsilon}_\nu U(t) \hat{\epsilon}_\mu = |\nu\rangle\langle\mu| \otimes \sum_{W(\mathcal{G})_{\nu\mu}} i^{-n} \int_0^t \int_0^{t_1} \dots \int_0^{t_{n-1}} U_\nu H_{\nu\leftarrow\eta_{n-1}} \dots U_{\eta_2} H_{\eta_2\leftarrow\mu} U_\mu dt_1 \dots dt_n \quad (\star)$$

Sum on strings of jumps of  $S'$  → Sum on **histories**  
**Walk-sum on  $\mathcal{G}$**

Continuous sums on jump times

$S'$  jumps,  $S$  affected by  $H_{\beta\leftarrow\alpha}$   
 $d_S \times d_S$  matrix

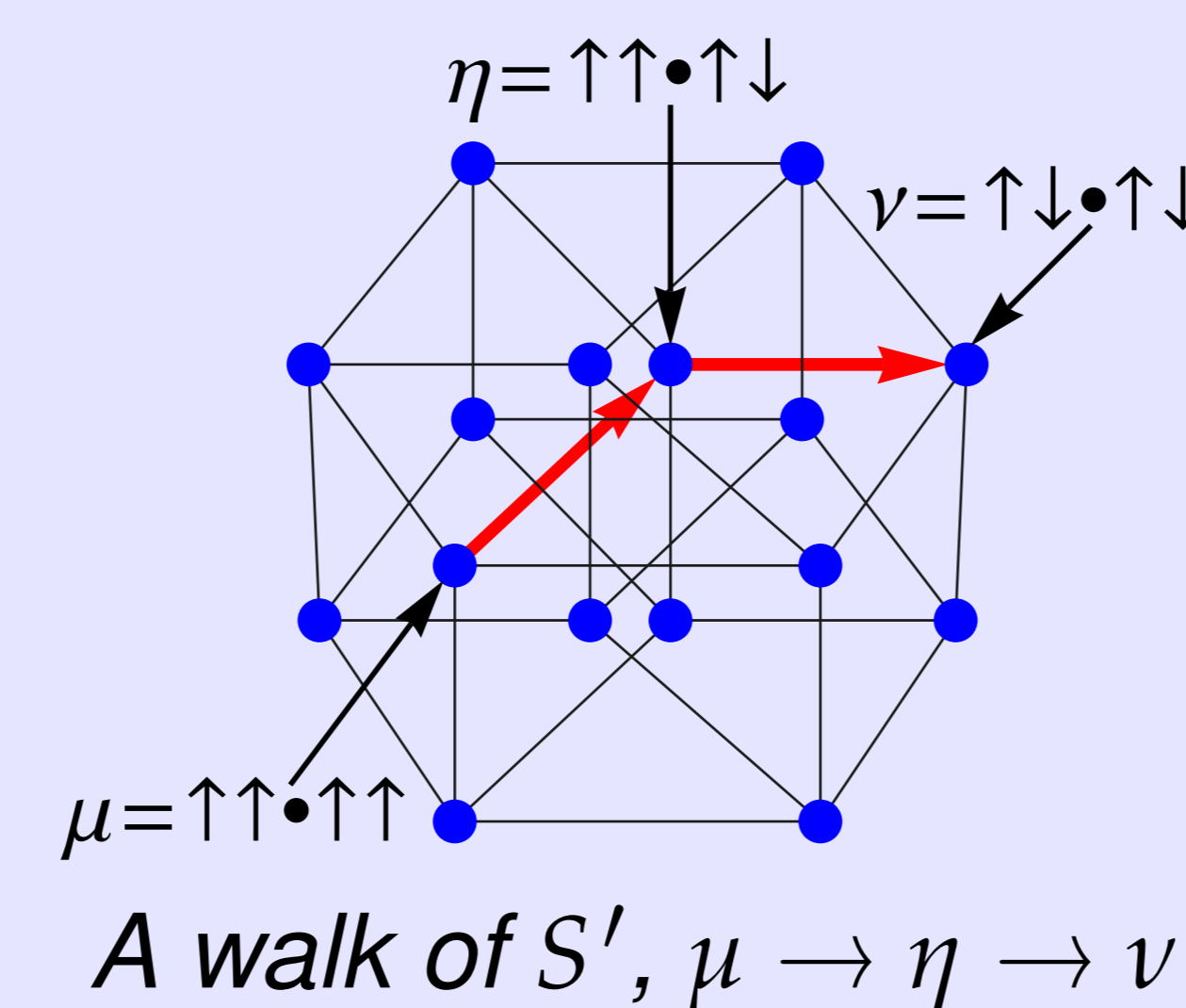
Evolution operator of  $S$  when  $S'$  is frozen in  $\mu$  from 0 to  $t_1$   
 $d_S \times d_S$  matrix

### Configuration space of $S'$ ...

... is a graph  $\mathcal{G}$  !

- Vertices = configurations
- Edges = allowed transitions

Ex: 5 spin-1/2 system,  
S central spin



## What walk-sum yields

We have obtained the *conditional* evolution operators  $U_{\nu\leftarrow\mu}$  defined by the walk-sum in  $(\star)$

- $U_{\nu\leftarrow\mu}$  is a  $d_S \times d_S$  piece of  $U$

$$U(t) = \sum_{\nu,\mu} \underbrace{|\nu\rangle\langle\mu|}_{S' \text{ from } \mu \text{ to } \nu} \otimes \underbrace{U_{\nu\leftarrow\mu}(t)}_{\text{Exact evolution of } S}$$

- Represents the evolution of  $S$ , knowing  $S'$  starts in  $\mu$  and finishes in  $\nu$
- Integrals are convolutions: Laplace domain → only + and  $\times$  of  $d_S \times d_S$  matrices required
- Effective matrices  $U_\nu$  and  $H_{\nu\leftarrow\mu}$  easy to compute !

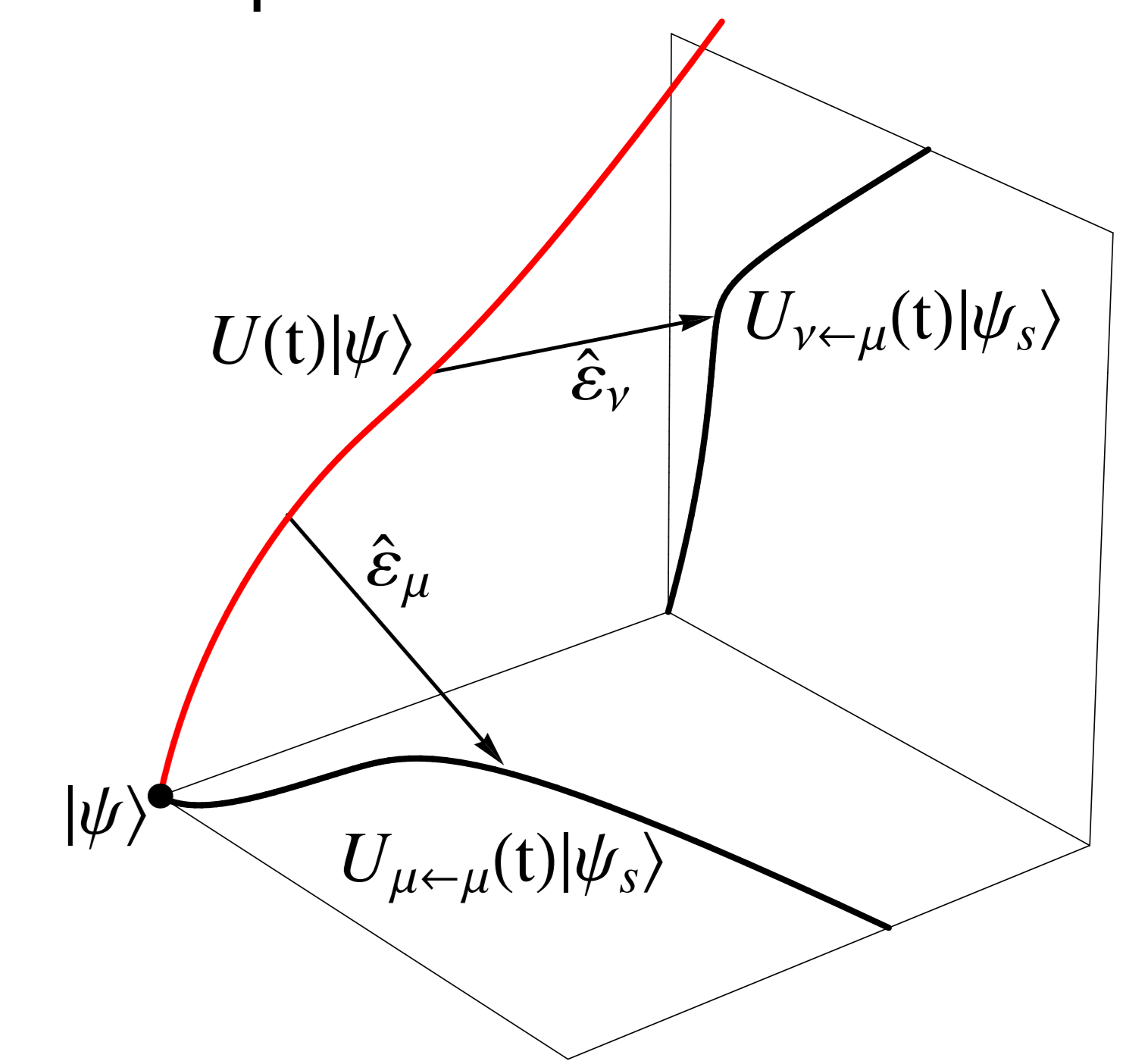
$$U_\nu = e^{-iH_\nu t} \quad \text{with} \quad \hat{\epsilon}_\nu H \hat{\epsilon}_\nu = \hat{\epsilon}_\nu \otimes H_\nu$$

$$\hat{\epsilon}_\nu H \hat{\epsilon}_\mu = |\nu\rangle\langle\mu| \otimes H_{\nu\leftarrow\mu}$$

→  $H_\nu$ : small effective Hamiltonian acting on  $S$  when  $S'$  frozen in  $\nu$

→  $H_\nu$  and  $H_{\nu\leftarrow\mu}$  are pieces of the full Hamiltonian  $H$

- $U_{\nu\leftarrow\mu}$  is expressed as a **sum of walks** = **sum of histories** of  $S'$
- Number of walks of length  $K$  on  $\mathcal{G}$  scales as **poly( $N$ )**



## Implementation, advantages and limitations

Unique procedure to follow: [arXiv:1204.5087](https://arxiv.org/abs/1204.5087)

- Divide the system into the two parts  $S$  and  $S'$
- Compute the effective  $H_\mu$  and  $H_{\nu\leftarrow\mu}$  required
- If possible, draw the graph  $\mathcal{G}$
- Decide which  $U_{\nu\leftarrow\mu}$  are of interest
- Compute these via the walk-sum

### Advantages

- Free to choose  $S$
- Can tackle 2D and 3D systems
- Observed to work best with **strong long-range** interactions
- Smooth monotonous convergence

### Limitations

- No rigorous criterion exists to decide when to implement walk-sum
- Walk-sum is a perturbation theory with the jumps of  $S'$  perturbing  $S$

### Rydberg excited Mott insulators

Model:

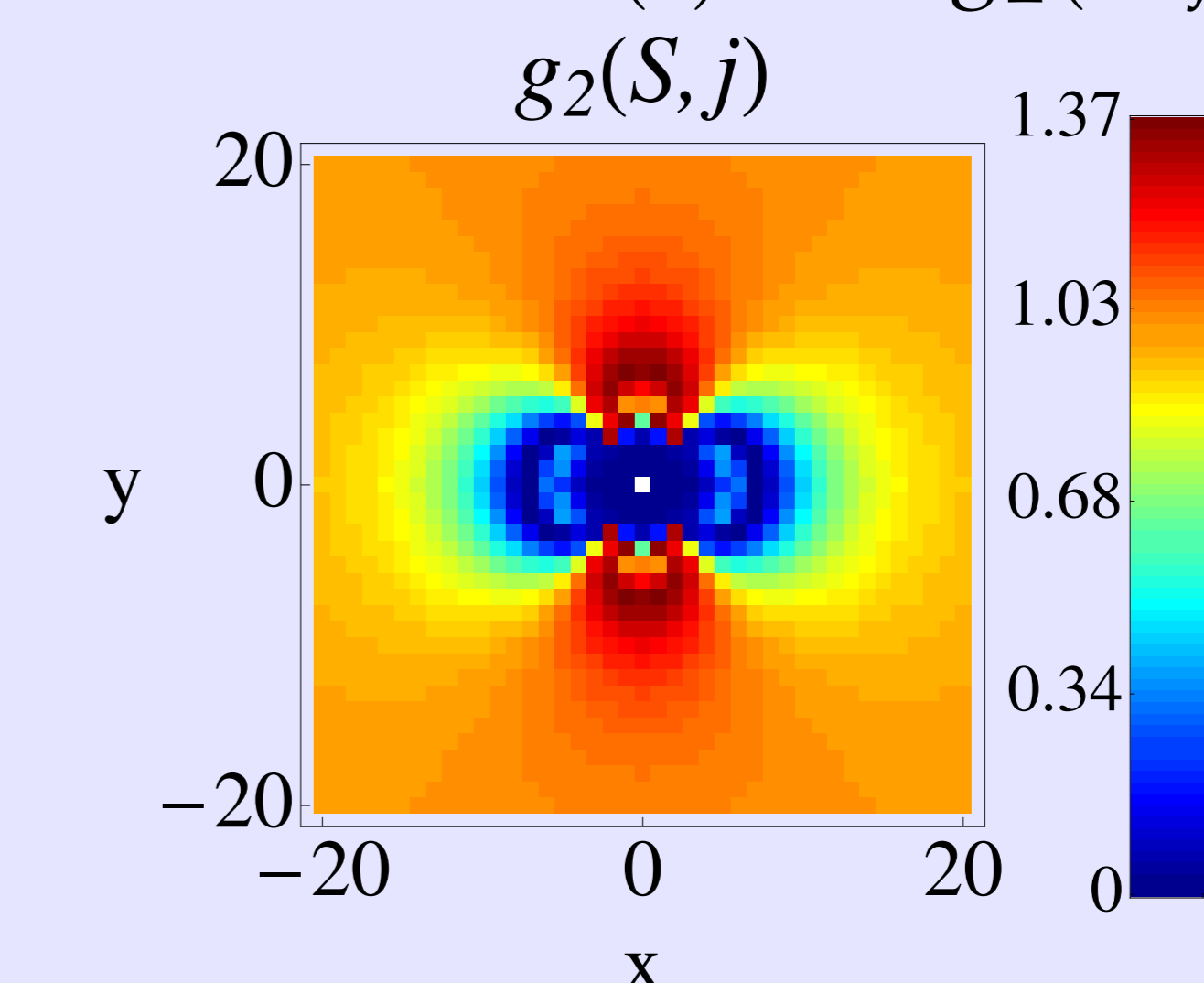
- Lattice of spin-1/2 + **strong long-range interaction**

$$H = \sum_i \left\{ \Delta |\uparrow_i\rangle\langle\uparrow_i| - \frac{1}{2} \Omega \sigma_x^i + \frac{1}{2} \sum_{j \neq i} \frac{A}{|i-j|^3} |\uparrow_i\uparrow_j\rangle\langle\uparrow_i\uparrow_j| \right\}$$

- Graph  $\mathcal{G} = N$ -hypercube
- Initial state  $|\downarrow_1 \dots \downarrow_N\rangle \equiv |0\rangle$

Density-density correlation function  $g_2(S, j)$

- $S$  central atom,
- Compute all  $U_{K\leftarrow 0}$  with  $K \leq 6$  excitations
- Reconstruct  $U(t)$  and  $g_2(S, j)$



- $A$ : dipole-dipole int.
- $\Delta = \Omega$ ,  $\pi$ -pulse
- 1 atom / pixel  
⇒  $N \simeq 1700$  atoms

More examples: [arXiv:1204.5087](https://arxiv.org/abs/1204.5087), [arXiv:1108.1177](https://arxiv.org/abs/1108.1177)

## From walk-sum to path-sum

Like walk-sum, **path-sum computes any piece**  $U_{v\leftarrow\mu}$  of  $U$ .

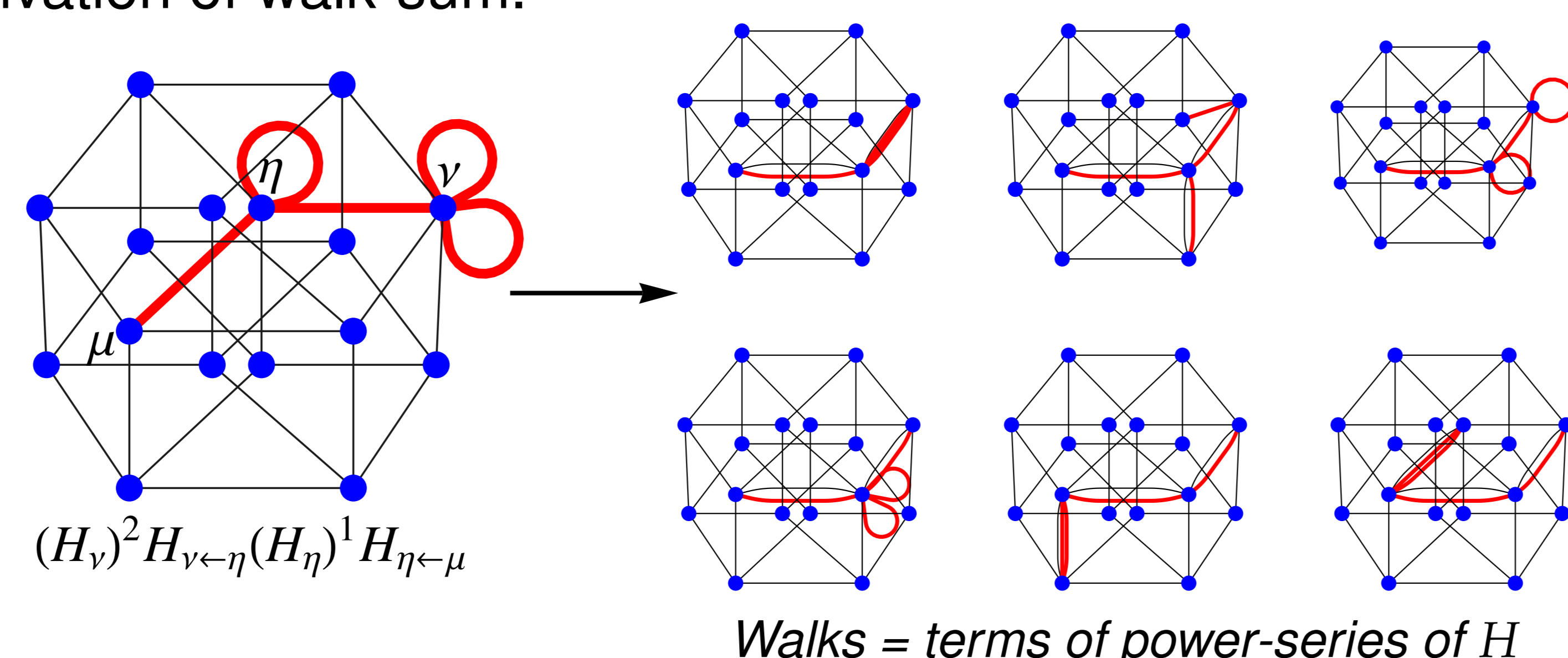
To understand path-sum, we take a new look on the derivation of walk-sum.

Graph  $\mathcal{G}_L$

- ▶ **Vertex:** configuration of  $S'$
- ▶ **Edge**  $\mu \rightarrow v$ : associate weight  $H_{v\leftarrow\mu}$
- ▶ **Walk weight:** product of edge weights

Remember  $H_\mu$  and  $H_{v\leftarrow\mu}$  = pieces of  $H$

→ Terms of power-series of  $H$  = walks on  $\mathcal{G}_L$



- ▶ The sum of walks differing only by loops can be calculated !

$$\begin{aligned}
 & \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \text{Diagram 4} + \dots = H_v H_{v\leftarrow\mu} e^{-iH_\mu t} \\
 & \frac{H_v H_{v\leftarrow\mu} (H_\mu)^0 (-i)^0}{0!} + \frac{H_v H_{v\leftarrow\mu} (H_\mu)^1 (-i)^1}{1!} + \frac{H_v H_{v\leftarrow\mu} (H_\mu)^2 (-i)^2}{2!} + \frac{H_v H_{v\leftarrow\mu} (H_\mu)^3 (-i)^3}{3!} + \dots
 \end{aligned}$$

The loops of  $\mathcal{G}_L$  are eliminated through exact resummations, thus yielding the **walk-sum** (★) on  $\mathcal{G}$

- ▶  $U_v = \exp[-iH_v t]$ : an effective vertex weight **representing** the sum of all loops
- ▶ **Path-sum sums all cycles** this way
  - Reduces the walk-sum to sum over acyclic walks: the simple paths
  - Effective vertex weight  $\tilde{U}_v$  representing the sum of all cycles

## Path-sum

- ▶ Path-sum gives the  $U_{v\leftarrow\mu}$  operators by summing over the **simple paths** of  $\mathcal{G}$ 
  - A simple path is forbidden to visit any vertex more than once
- ▶ Path-sum is non-perturbative
- ▶ Path-sum yields a generalized Dyson equation
  - Path-sum also gives an explicit formula for the self-energy !

Conclusion:

Features	Walk-sum	Path-sum
Physical derivation	✓	×
Perturbative	✓	×
Valid in 2D and 3D	✓	✓
Valid for any choice of $S$	✓	✓
Valid for any Hamiltonian	✓	✓
Valid for any matrix function	×	✓
Criterion to decide practicality	×	×
Smooth monotonous convergence	✓	?
Systems with continuous degrees of freedom	Path-Integrals	✓

**Walk-sum:**

Works best for systems with strong long-range interactions. A discrete equivalent to Feynman path-integrals.

**Path-sum:**

Generalized Dyson equation valid for any  $S$ , explicitly gives the self-energy. Works for all matrices, all functions.

## The path-sum formula

- ▶ Gives  $U_{v\leftarrow\mu}$  in the Laplace domain

$$U_{v\leftarrow\mu}(t) = \mathcal{L}^{-1} \left\{ \sum i^{-n} \tilde{U}_{\mathcal{G}\setminus\{\mu, \dots, \eta_n\}}[v] \left[ H_{v\leftarrow\eta_n} \cdots \tilde{U}_{\mathcal{G}\setminus\{\mu\}}[\eta_2] H_{\eta_2\leftarrow\mu} \tilde{U}_{\mathcal{G}}[\mu] \right] \right\} (t)$$

Sum over the **simple paths**

Effective evolution operator of  $S$  while  $S'$  undergoes arbitrary cyclic evolution off  $\mu$   
 Represents the sum of all cycles off  $\mu$  on  $\mathcal{G}$

$$\tilde{U}_{\mathcal{G}}[\mu] = \left[ s\mathcal{I} - H_\mu - \sum i^{-m} H_{\mu\leftarrow\lambda_m} \tilde{U}_{\mathcal{G}\setminus\{\mu, \dots, \lambda_{m-1}\}}[\lambda_m] \cdots \tilde{U}_{\mathcal{G}\setminus\{\mu\}}[\lambda_2] H_{\lambda_2\leftarrow\mu} \right]^{-1} \iff G = [G_0^{-1} - \Sigma]^{-1}$$

Sum over the **simple cycles**

Progressively smaller configuration space

Generalized Dyson equation

- ▶ Valid for any  $S$ , discrete degrees of freedom
- ▶  $G$  dressed propagator:  $\tilde{U}_\mu$  evolves  $S$  exactly
- ▶  $G_0$  bare propagator:  $[s\mathcal{I} - H_\mu]^{-1}$  evolves  $S$  with  $S'$  static
- ▶  $\Sigma$  self energy: sum over the simple cycles
  - Finite explicit formula for  $\Sigma$

Derivation: [arXiv:1202.5523](https://arxiv.org/abs/1202.5523)

## How do we use path-sum?

Demonstrative example: two atoms

- ▶ Each atom has 2 levels:  $g \equiv \bullet$  and  $e \equiv \blacktriangleright$
- ▶  $S =$  atom 1
  - $S'$  atom 2  $\Rightarrow \mathcal{G}$  has 2 vertices
- ▶ What is  $U_{e\leftarrow g}$ ? **one simple path** ...

$$U_{e\leftarrow g} = \mathcal{L}^{-1} \left\{ \tilde{U}_{\mathcal{G}\setminus\{g\}}[e] H_{e\leftarrow g} \tilde{U}_{\mathcal{G}}[g] \right\}^{-1}$$

... and a sum on the simple cycles

$$\tilde{U}_{\mathcal{G}}[g] = [s\mathcal{I} - H_g - H_{g\leftarrow e} \tilde{U}_{\mathcal{G}\setminus\{g\}}[e] H_{e\leftarrow g}]^{-1}$$

$$\tilde{U}_{\mathcal{G}\setminus\{g\}}[e] = [s\mathcal{I} - H_e]^{-1}$$

General matrix functions: [arXiv:1112.1588](https://arxiv.org/abs/1112.1588)

