

The Theory of Operator-Lattices

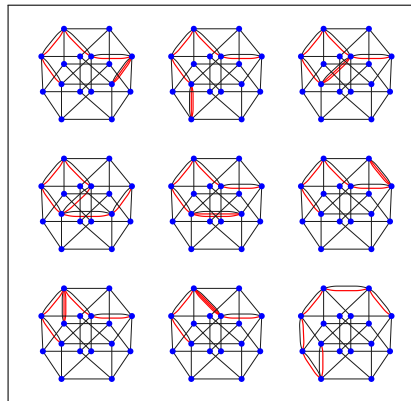
A New Analytical Tool for Quantum Many-Body Physics

P.L. Giscard



September 17, 2010

- 1 Original Motivation
- 2 Conditional-Evolution Operators
- 3 Graphical Representation
- 4 Rydberg Atoms
- 5 Summary & Future Research



Original Motivation

GOAL

Efficiently find the time dynamics of a many-body system

- System with N interacting 2-level atoms

$$\mathcal{H} \Rightarrow 2^N \times 2^N \gg 1 \Rightarrow \mathbf{Hard}$$

- Same system, all atoms frozen except 1

$$\mathcal{H} \Rightarrow 2 \times 2 \Rightarrow \mathbf{Easy}$$

Can we express the true evolution as \mathcal{F} (frozen situations) ?

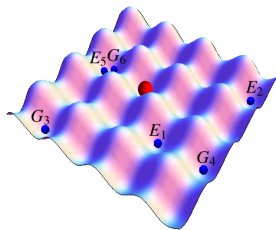
Operator-Lattices

Freezing atoms = put + keep in some state $|\psi_\mu\rangle$

- Operator-lattice \Rightarrow lattice of projectors

$$\varepsilon_\mu = P_\mu \otimes \mathcal{I}_{atom}$$

$$P_\mu = |\psi_\mu\rangle\langle\psi_\mu|$$



$$\varepsilon_\mu = E_1 E_2 G_3 G_4 E_5 G_6 \otimes \mathcal{I}_{red}$$

Freezing the System

- Introduce closure relation $U(t) = \lim_{\delta t \rightarrow 0} \prod_n^N \left(\sum_{\mu} \varepsilon_{\mu} \right) \delta U$
 and *expand*

$$U = \lim_{\delta t \rightarrow 0} \varepsilon_{\mu} \delta U \varepsilon_{\mu} \delta U \dots \varepsilon_{\mu} \delta U \varepsilon_{\mu} +$$

$$\lim_{\delta t \rightarrow 0} \varepsilon_{\mu} \delta U \varepsilon_{\mu} \delta U \dots \varepsilon_{\mu} \delta U \varepsilon_{\nu} \dots \varepsilon_{\nu} \delta U \varepsilon_{\nu} + \dots$$

Zeno measurement of ε_{μ}

Zeno measurement of ε_{ν}
 + jump +
 Zeno measurement of ε_{μ}

- Full expansion = *all* jumps + *all* jumping times

Conditional-Evolution Operators

$$U_{\nu \rightarrow \mu}(t) = \sum_n (i\hbar)^{-n} \sum_{P_n(G)} \int_0^t \int_0^{t_n} \dots \int_0^{t_2} e^{-iH_\mu(t-t_n)/\hbar} H_{\eta \rightarrow \mu} \dots H_{\nu \rightarrow \chi} e^{-iH_\nu t_1/\hbar} dt_1 \dots dt_n$$

Number of jumps

Strings of n
consecutive jumps

Times of the jumps

- H_μ effective Hamiltonian when system in ε_μ , 2×2 matrix

$$\varepsilon_\mu \mathcal{H} \varepsilon_\mu = H_\mu \otimes \varepsilon_\mu$$

- $H_{\eta \rightarrow \mu}$ how a jump from ε_η to ε_μ affects the atom, 2×2 matrix

$$\varepsilon_\mu \mathcal{H} \varepsilon_\eta = H_{\eta \rightarrow \mu} \otimes |\varepsilon_\mu\rangle \langle \varepsilon_\eta|$$

Conditional-Evolution Operators

Conditional-evolution operators have nice properties !

- $U_{\nu \rightarrow \mu}(t) = \varepsilon_{\mu} U(t) \varepsilon_{\nu}$
- They evolve quantum conditional expectation values
$$\langle A / \varepsilon_{\mu} \rangle = \langle \psi_{\nu} | U_{\nu \rightarrow \mu}^{\dagger} A U_{\nu \rightarrow \mu} | \psi_{\nu} \rangle$$
- Fourier domain $\tilde{U}_{\nu \rightarrow \mu}(\omega)$ only \times and $+$ of **universal 2×2** matrices
- Can be interpreted as interference of information TSVF

Future \longrightarrow *Past* & *Past* \longrightarrow *Future*

The Simple Graph

*Is it **practical** for calculations ?*

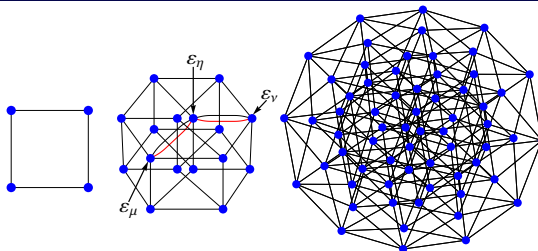
$$U_{\nu \rightarrow \mu}(t) \leftarrow \sum_n \sum P_n(G)$$

- Huge number of jumps ?
- Find all the possible strings of jumps ? some $H_{\eta \rightarrow \chi} = 0$

Simple graph G

- For each operator-lattice ε_μ place a vertex
- If $H_{\eta \rightarrow \mu} \neq 0$ draw an edge between ε_μ and ε_η

The Simple Graph



Graphs G for 1D, 2D and 3D XYZ-Hamiltonians

- String of jumps = Path on G
- Number of strings = Number of paths on G
- N.n Hamiltonians reduce to small graphs

Path-integrals on G

The Speed-Up

*Are calculations **efficient** ?*

- Number of operations

$$U(t) \propto 2^{3N} K$$
$$U_{\nu \rightarrow \mu}(t) \propto N^K$$

$U_{\nu \rightarrow \mu}$ generates a small piece of the full $U(t)$

- Operations / elements of $U(t)$ generated

$$U(t) \propto 2^N$$
$$U_{\nu \rightarrow \mu}(t) \propto N^K$$

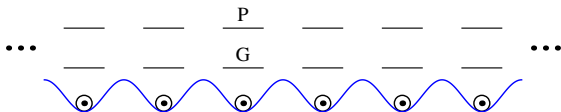
Exponential speed-up

The Underlying Mathematical Theory

Function of matrices = Paths on graphs

- Extension to any matrix-function
- Derivation all the results independently of physics
- Exponential speed-up on $\mathbf{Exp}(\mathcal{M})$ and \mathcal{M}^{-1}

Rydberg Atoms in Deep Optical Lattices



$\langle P \rangle$? *Distribution of Rydberg excitations ?*

$$\mathcal{H} = \sum_i \mathcal{H}_{\text{Laser}} + \sum_{i,j} \frac{A}{R^3} P^i P^j$$

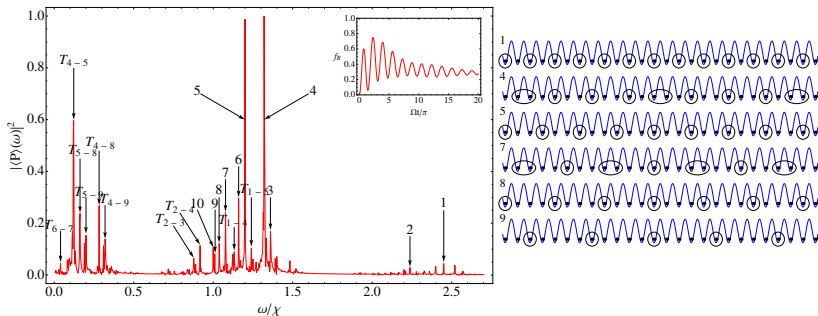


Strong anisotropic dipole-dipole interaction

Blockade, antiblockade \Rightarrow ~~Mean Field~~

- MF-based $\frac{A}{R^3} P^i P^j \simeq \frac{A}{R^3} P^i \langle P \rangle g_2$
- Using operator-lattices $\langle P \rangle = \mathcal{F}(\langle P^i / P^j \rangle, \langle P^i / G^j \rangle)$

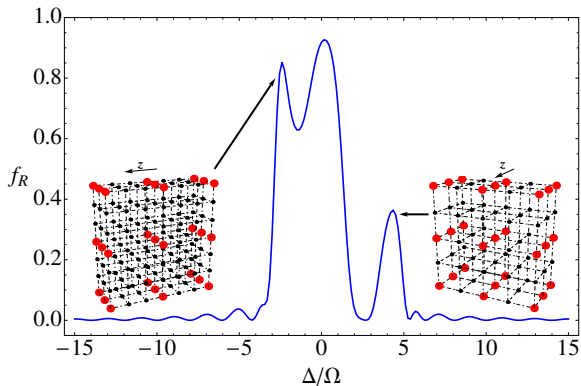
Results



- Lowest order
- n -atom correlation functions
- Rydberg fraction

- Lattice coverings
- Quantum phase-transition

Results



- 2D and 3D
- Selective population of lattice-coverings

Summary

"Path integrals suffer most grievously from a serious defect. They do not permit a discussion of spin operators or other such operators in a simple and lucid way."

R. Feynman

Main results

- Path-integrals (sums) on discrete spaces
- System of N particles with n_ℓ levels.

Time-dynamics \Rightarrow \times and $+$ of $n_\ell \times n_\ell$ matrices

- Large speed-up in computations
- General theory of matrix functions, *independent of physics*

Future Research

Perspectives

- Systematic application to physical systems
- Integration of damping
- Link with TSVF
- Recovering usual path-integrals (continuous G)
- Efficient computation of the matrix-log